

Physics 2102

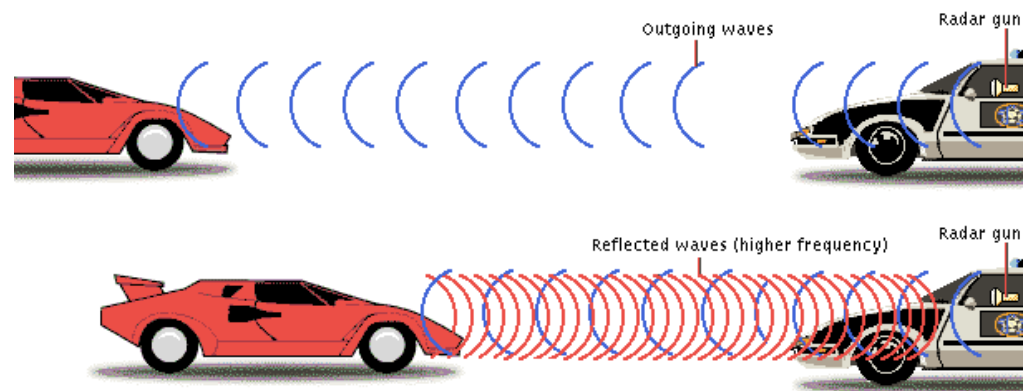
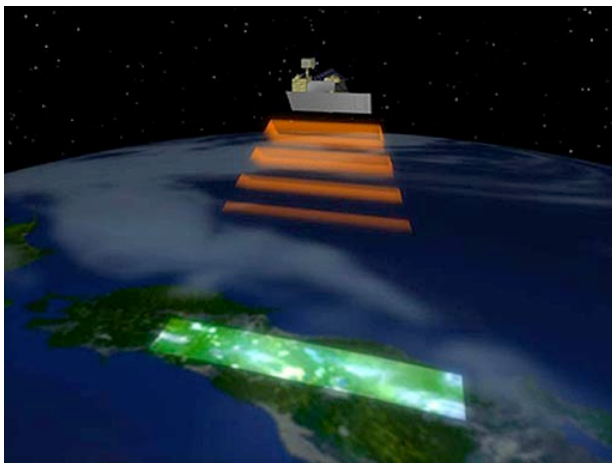
Christian Buth



Lecture 41

Diffraction 2

04/29/2009



Review

- Interference only for coherent light, i.e., with a **phase relationship** that is time independent
- **Intensity** in double-slit interference:

$$I = 4I_0 \cos^2 \frac{1}{2} \phi$$

$$\phi = \frac{2\pi d}{\lambda} \sin \theta$$

- Use Huygens' Principle to find positions of **diffraction minima** of a single slit by subdividing the aperture

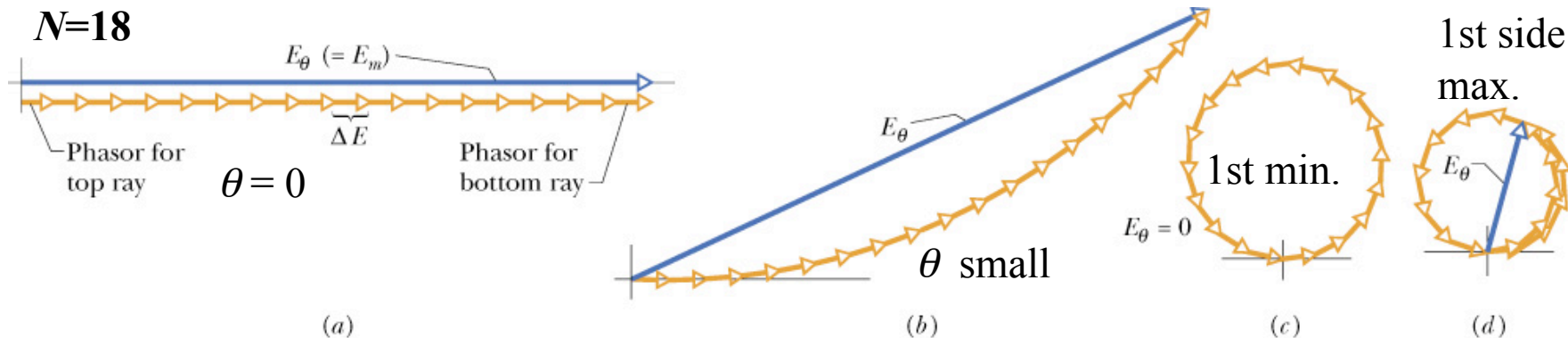
$$a \sin \theta = m\lambda, \text{ for } m = 1, 2, 3 \dots \text{ (minima-dark fringes)}$$

Intensity in Single-Slit Diffraction, Qualitatively

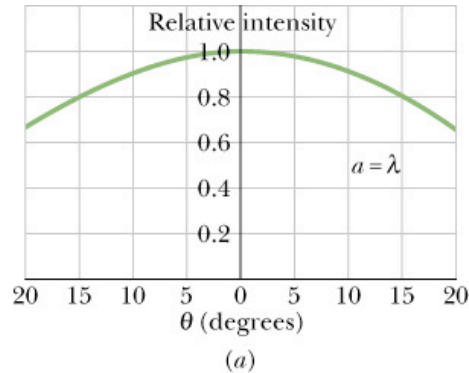
To obtain the locations of the minima, the slit was equally divided into N zones, each with width Δx . Each zone acts as a source of Huygens wavelets. Now these zones can be superimposed at the screen to obtain the intensity as a function of θ , the angle to the central axis.

To find the net electric field E_θ (intensity $\propto E_\theta^2$) at point P on the screen, we need the phase relationships among the wavelets arriving from different zones:

$$\left(\begin{array}{c} \text{phase} \\ \text{difference} \end{array} \right) = \left(\frac{2\pi}{\lambda} \right) \left(\begin{array}{c} \text{path length} \\ \text{difference} \end{array} \right) \quad \Delta\phi = \left(\frac{2\pi}{\lambda} \right) (\Delta x \sin \theta)$$



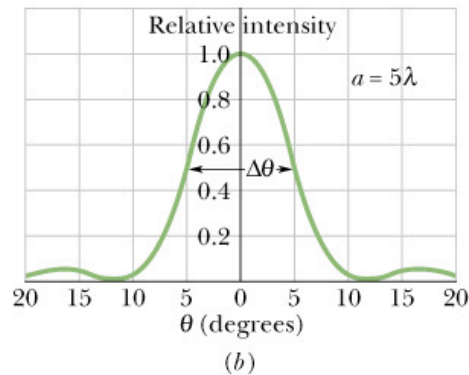
Intensity in Single-Slit Diffraction, Quantitatively



Here we will show that the intensity at the screen due to a single slit is:

$$I(\theta) = I_m \left(\frac{\sin \alpha}{\alpha} \right)^2 \quad (36-5)$$

$$\text{where } \alpha = \frac{1}{2} \phi = \frac{\pi a}{\lambda} \sin \theta \quad (36-6)$$



In Eq. 36-5, minima occur when:

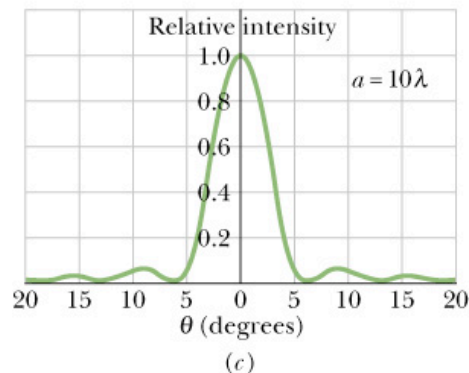
$$\alpha = m\pi, \quad \text{for } m = 1, 2, 3, \dots$$

If we put this into Eq. 36-6 we find:

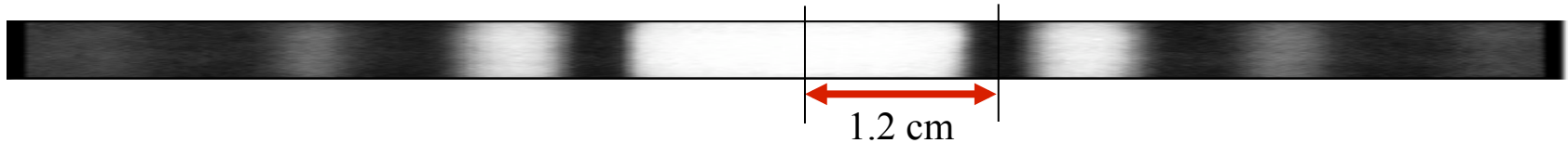
$$m\pi = \frac{\pi a}{\lambda} \sin \theta, \quad \text{for } m = 1, 2, 3, \dots$$

$$\text{or } a \sin \theta = m\lambda, \quad \text{for } m = 1, 2, 3, \dots$$

(minima-dark fringes)



Diffraction of a laser through a slit



Light from a helium-neon laser ($\lambda = 633 \text{ nm}$) passes through a narrow slit and is seen on a screen 2.0 m behind the slit. The first minimum of the diffraction pattern is observed to be located 1.2 cm from the central maximum.

How wide is the slit?

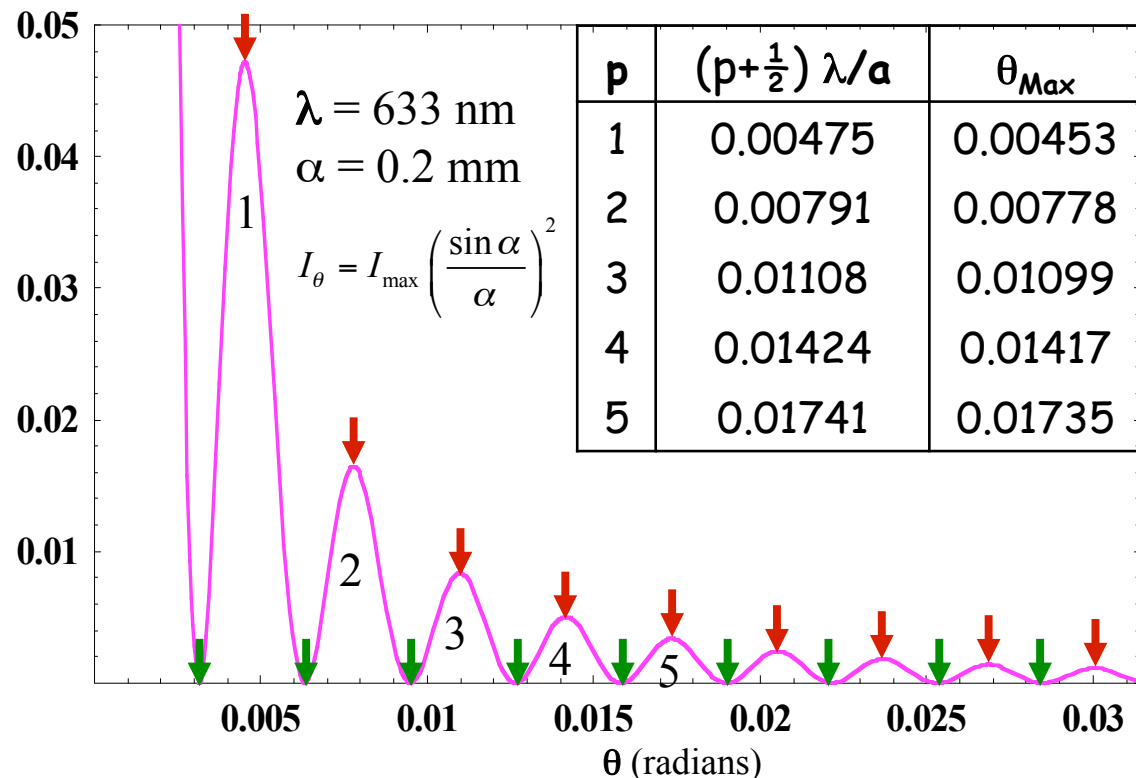
$$\theta_1 = \frac{y_1}{L} = \frac{(0.012 \text{ m})}{(2.00 \text{ m})} = 0.0060 \text{ rad}$$

$$a = \frac{\lambda}{\sin \theta_1} \cong \frac{\lambda}{\theta_1} = \frac{(6.33 \times 10^{-7} \text{ m})}{(6.00 \times 10^{-3} \text{ rad})} = 1.06 \times 10^{-4} \text{ m} = 0.106 \text{ mm}$$

Angles of the Secondary Maxima

The diffraction minima are *precisely* at the angles where
 $\sin \theta = p \lambda/a$ and $\alpha = p\pi$
 (so that $\sin \alpha = 0$).

However, the diffraction maxima are *not quite* at the angles where
 $\sin \theta = (p+1/2) \lambda/a$
 and $\alpha = (p+1/2)\pi$
 (so that $|\sin \alpha| = 1$).



To find the maxima, one must look *near* $\sin \theta = (p+1/2) \lambda/a$, for places where the slope of the diffraction pattern goes to zero, i.e., where $d[(\sin \alpha/\alpha)^2]/d\theta = 0$. This is a transcendental equation that must be solved numerically. The table gives the θ_{Max} solutions. Note that $\theta_{\text{Max}} < (p+1/2) \lambda/a$.

Diffraction by a Circular Aperture

Distant point
source, e.g., star

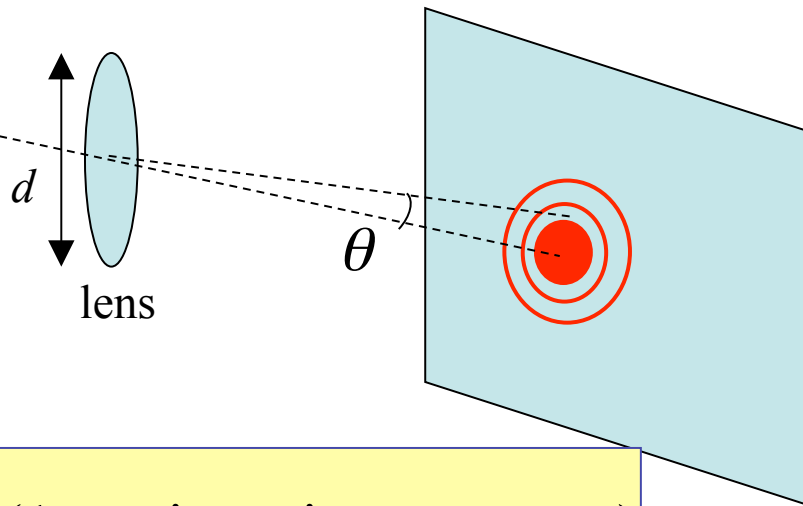
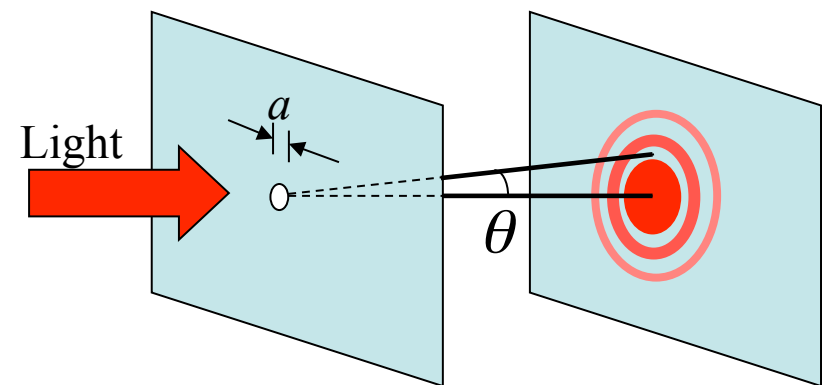
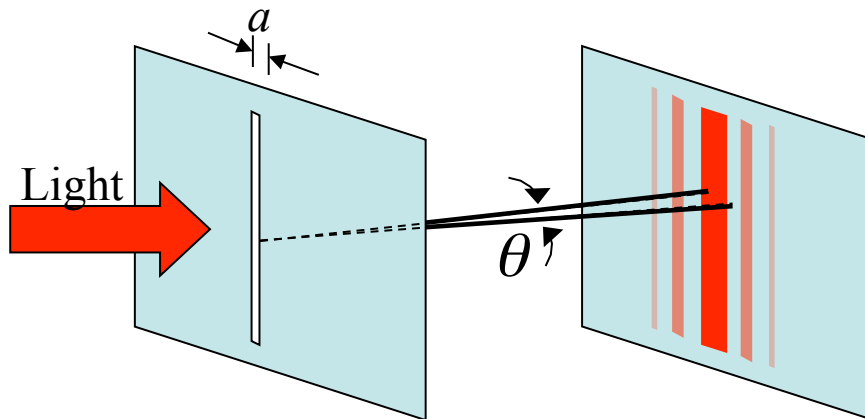


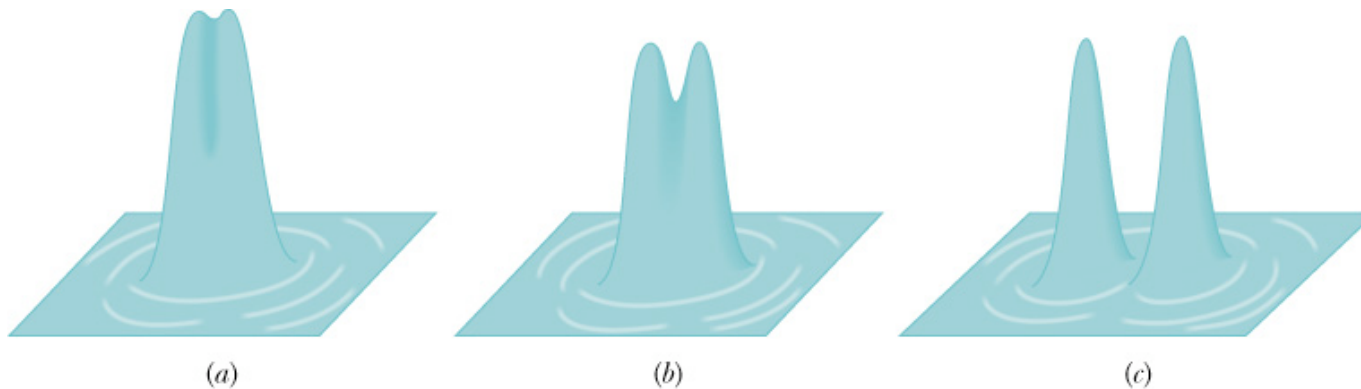
Image is not a point, as expected from geometrical optics! Diffraction is responsible for this image pattern.

$$\sin \theta = 1.22 \frac{\lambda}{d} \text{ (1st min.- circ. aperture)}$$



Resolvability

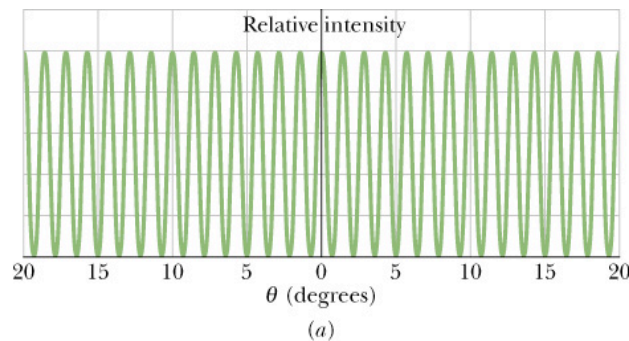
Rayleigh's Criterion: Two point sources are barely resolvable if their angular separation θ_R results in the central maximum of the diffraction pattern of one source's image centered on the first minimum of the diffraction pattern of the other source's image



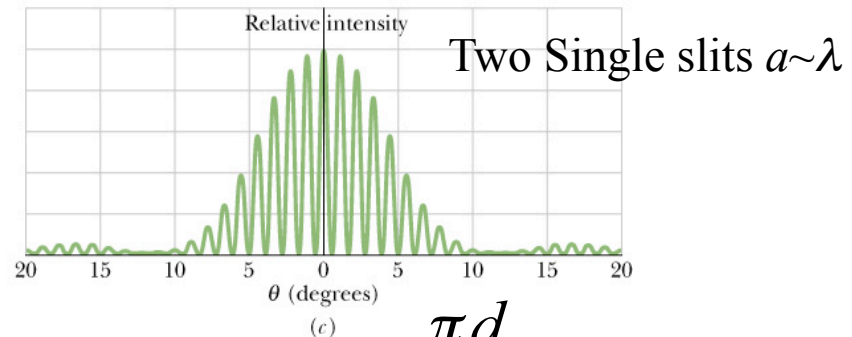
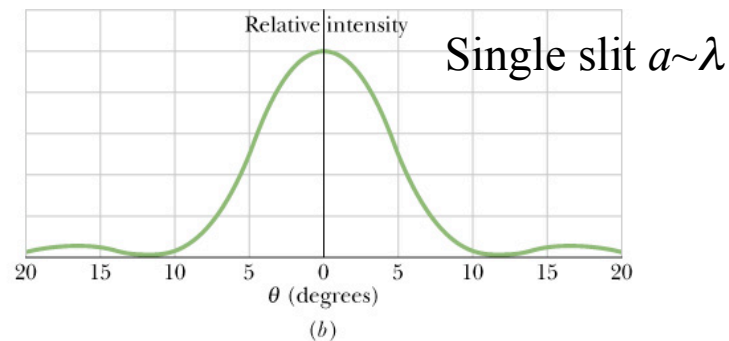
$$\theta_R = \sin^{-1} \left(1.22 \frac{\lambda}{d} \right)^{\theta_R \text{ small}} \approx 1.22 \frac{\lambda}{d} \quad (\text{Rayleigh's criterion})$$

Diffraction by a Double Slit

In the double-slit experiment, we assumed that the slit width $a \ll \lambda$. What if this is not the case?



Two vanishingly narrow slits $a \ll \lambda$



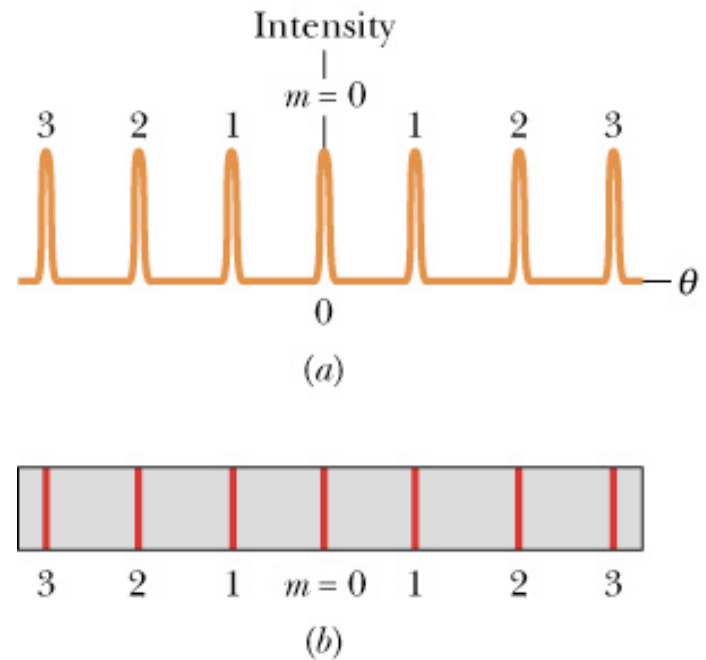
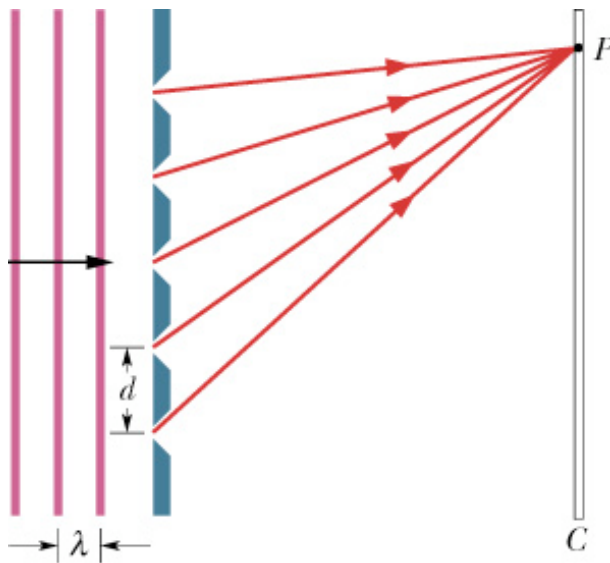
$$I(\theta) = I_m \left(\cos^2 \beta \right) \left(\frac{\sin \alpha}{\alpha} \right)^2 \quad (\text{double slit})$$

$$\beta = \frac{\pi d}{\lambda} \sin \theta$$

$$\alpha = \frac{\pi a}{\lambda} \sin \theta$$

Diffraction Gratings

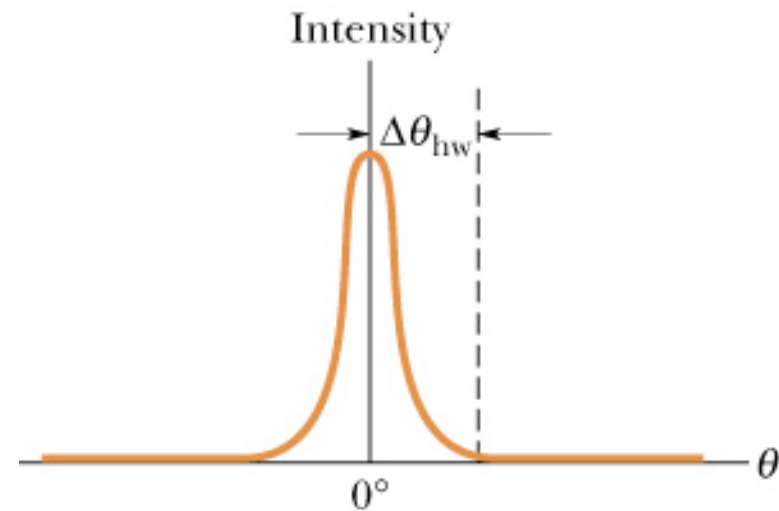
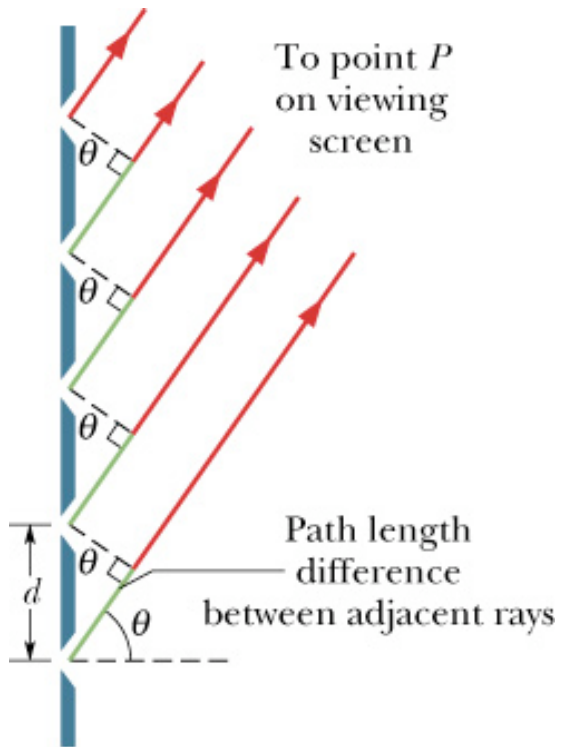
A device with N slits (rulings) can be used to manipulate light, such as separate different wavelengths of light that are contained in a single beam. How does a diffraction grating affect monochromatic light?



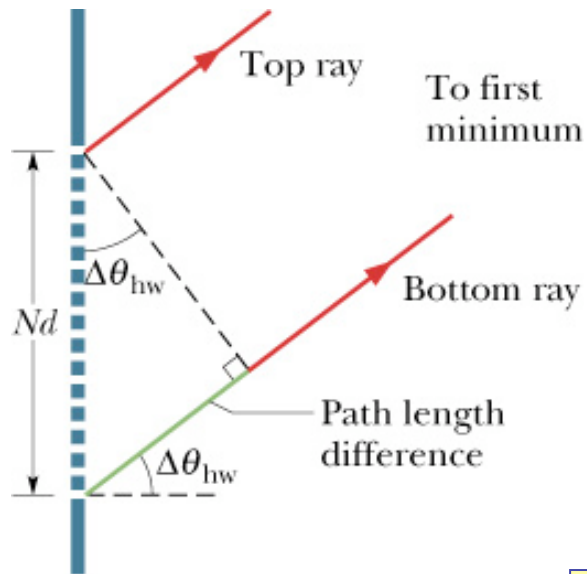
$$d \sin \theta = m\lambda \quad \text{for } m = 0, 1, 2, \dots \quad (\text{maxima-lines})$$

Width of Lines

The ability of the diffraction grating to resolve (separate) different wavelengths depends on the width of the lines (maxima).



Width of Lines, cont'd



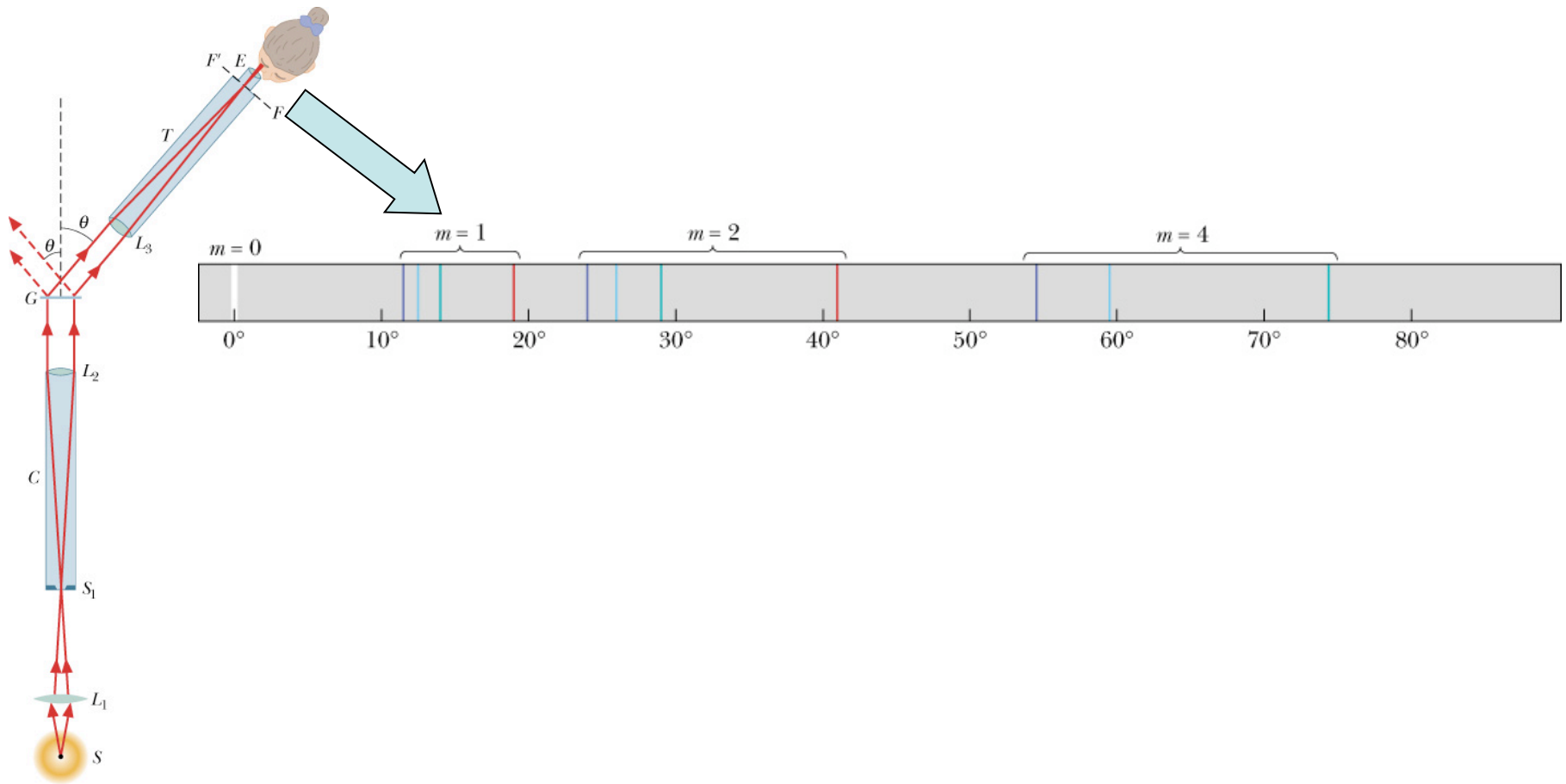
$$Nd \sin \Delta\theta_{hw} = \lambda, \quad \sin \Delta\theta_{hw} \approx \Delta\theta_{hw}$$

$$\Delta\theta_{hw} = \frac{\lambda}{Nd} \quad (\text{half width of central line})$$

$$\Delta\theta_{hw} = \frac{\lambda}{Nd \cos \theta} \quad (\text{half width of line at } \theta)$$

Grating Spectroscope

Separates different wavelengths (colors) of light into distinct diffraction lines



Summary

- To predict the interference pattern of a multi-slit system, we must combine interference and diffraction effects.
- **Rayleigh's Criterion** for separability of two points
- Intensity in single-slit diffraction:

$$I(\theta) = I_m \left(\frac{\sin \alpha}{\alpha} \right)^2 \quad \text{where } \alpha = \frac{1}{2} \phi = \frac{\pi a}{\lambda} \sin \theta \quad (36-6)$$

- Double-slit diffraction:

$$I(\theta) = I_m \left(\cos^2 \beta \right) \left(\frac{\sin \alpha}{\alpha} \right)^2 \quad (\text{double slit})$$